

# Design Procedure for Interdigital Waveguide Bandpass Filters based on Flat Metallic Strips

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## Key Points:

- A new topology for interdigital bandpass filters based on flat metallic strips is proposed
- A systematic procedure based on the segmentation of the component is presented for the efficient design of the proposed interdigital bandpass filters
- A power handling capability study of the designed interdigital filters is performed

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## Abstract

In this work, we describe a new topology for interdigital bandpass filters implemented in waveguide technology. Compared to classical coaxial resonators, which are typically based on cylindrical metallic posts, the resonators that we propose are based on rectangular boxed cavities loaded with flat metallic strips of a finite thickness. The main advantages of the proposed topology are low-cost practical realization and manufacture simplicity. The design process, which is based on the classical circuit-based approach combined with an efficient segmentation of the component, uses the well-known Aggressive Space Mapping technique, with the aim of reducing the overall computation effort. In order to validate the new topology and the design procedure, interdigital bandpass filters of orders 3 and 5 operating at the S-band have been designed. For validation purposes, the results obtained are successfully compared to simulations provided by two different full-wave electromagnetic tools (i.e., Ansys HFSS and CST Studio Suite). Furthermore, a multipactor study has also been performed, with the aim of exploring the power handling capability of the proposed filters in view of possible space applications. Finally, the design of classical interdigital filters based on cylindrical posts is also discussed, with the aim of comparing both topologies regarding compactness, as well as in terms of their electrical and high-power performance.

## 1 Introduction

Bandpass filters based on coaxial resonators, such as filters in combline or interdigital configuration implemented in waveguide technology, are commonly used in multiple applications, including space and wireless communications systems (Wang & Yu, 2009), (García et al., 2018), (Macchiarella et al., 2018), (Bastioli et al., 2018), (Vallerotonda et al., 2019), (Tamiazzo, 2021), (Zeng et al., 2022), (Yan et al., 2024). Classically, resonators based on rectangular cavities loaded with metallic posts of cylindrical (J. Li et al., 2018), (Boni et al., 2023) or rectangular/square geometry (Arndt & Brandt, 2002), (Kong et al., 2023) have been used in this type of coaxial resonator filters. Other more complex topologies have also been investigated in the past, as the interdigital bandpass filter designed in (Venter et al., 2020), where resonator posts with a triangular cross-section were proposed. Generally, coaxial resonators are coupled by using irises of rectangular geometry, whose position in the structure determines the type of coupling (electric or magnetic) (Jamshidi-Zarmehri et al., 2023), (Zeng et al., 2023). On the other hand, inter-resonator coupling can also be controlled by designing the separation between the posts, without the need to insert irises in the component (Z. Li et al., 2023), (Macchiarella et al., 2023).

Regarding the state-of-the-art of interdigital bandpass filters based on cylindrical posts, an in-line topology was presented in (Ikram & Rana, 2006), where a fifth-order Chebyshev interdigital bandpass filter operating at C-band was designed. Moreover, a manifold cavity diplexer operating at L-band was designed in (Packiaraj et al., 2005), while a low-loss triplexer that used a combination of a microstrip star junction and three interdigital bandpass filters was proposed in (Munina & Turalchuk, 2015). More recently, a Ku-band 8th-order interdigital filter with an in-line configuration was presented in (J. Li et al., 2018).

On the other hand, resonators based on rectangular or square posts have also been widely used in interdigital waveguide filters. For instance, an interdigital cavity filter operating at Ku-band is designed in (Wu et al., 2019), where two transmission zeros are realized in the frequency response by considering additional resonators outside of the ports of the filter. A further example of an interdigital filter using posts with a rectangular geometry can be found in (Jiang et al., 2019), where an eighth-order filter with an improved stop-band suppression is designed. Moreover, a wideband waveguide diplexer based on interdigital filters for mobile base stations is proposed in (Kobrin et al., 2019). Another

65 interesting contribution was presented in (Anand, 2021), where a folded configuration  
 66 for interdigital filters was proposed in order to implement cross-couplings between the  
 67 resonators. Additive manufacturing techniques have also been used for designing inter-  
 68 digital filters for radar applications (Rodriguez-Morales et al., 2021), and satellite com-  
 69 munication systems (Venter et al., 2020). More recently, a miniaturized interdigital cavi-  
 70 ty filter using loading capacitors is investigated in (Kong et al., 2023). In general, the  
 71 design process for the bandpass filters presented in the contributions discussed above,  
 72 is based on the classical approach proposed in (Cameron et al., 2018), so, regardless of  
 73 the filter order, all the relevant dimensions of the component are optimized at the same  
 74 time, thus requiring a great computational effort to complete the design process.

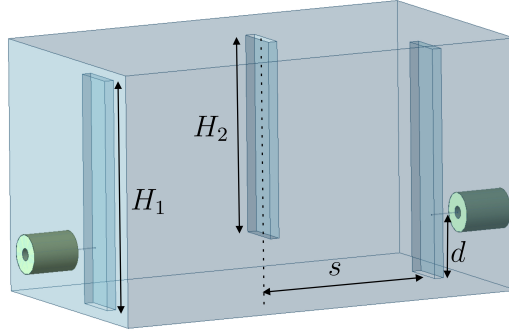
75 In this work, a new topology of waveguide bandpass filters based on coaxial reso-  
 76 nators implemented in an interdigital configuration is presented. The main novelty of  
 77 the proposed in-line filter consists in the use of cavities loaded with planar metallic strips  
 78 of a finite thickness, as shown in Fig. 1, where a third-order interdigital bandpass filter  
 79 is depicted. A major advantage of this new topology lies in the fact that the filters could  
 80 be manufactured using low-cost and more simple realization techniques. For instance,  
 81 the whole structure can be implemented through the cascaded connection of simpler rect-  
 82 angular waveguide sections (including the flat strips), each one manufactured with a CNC  
 83 (computer numerical control) milling technique. The design procedure that we propose  
 84 combines the use of the classical circuit-based approach (Cameron et al., 2018) with a  
 85 strategy based on a segmentation of the structure (San-Blas et al., 2021), (Jamshidi-Zarmehri  
 86 et al., 2023).

87 In the suggested design procedure, input and output resonators are first designed  
 88 to achieve the desired external quality factor. After that, the coupling between adjacent  
 89 resonators is addressed, with the aim of achieving a set of initial values for the physi-  
 90 cal dimensions of the structure. To continue, and specially for high-order filters, a sys-  
 91 tematic design procedure based on the segmentation of the whole component is proposed.  
 92 In this procedure, the different resonators of the filter are added one after the other, thus  
 93 dividing the design process in more simple steps, while also reducing the number of vari-  
 94 ables to be optimized in each stage. Equivalent circuit models are also important in this  
 95 approach, since their electrical responses will be used as target curves in all optimiza-  
 96 tion stages of the design procedure. Furthermore, the Aggressive Space Mapping (ASM)  
 97 technique (Bandler et al., 1995) has also been used in this work, so that most of the op-  
 98 timizations performed during the design process are carried out very efficiently in a low-  
 99 precision simulation space (the coarse model).

100 In order to validate the new topology that we propose, two interdigital bandpass  
 101 filters operating at S-band have been designed. The correct electrical performance of the  
 102 filters is successfully validated using two different electromagnetic (EM) simulation tools  
 103 (Ansys HFSS, and CST Studio Suite). In addition, a multipactor study (Berenguer et  
 104 al., 2019) is also performed for both filters, in order to evaluate the high-power perfor-  
 105 mance of this new topology under high-vacuum conditions. Furthermore, the design of  
 106 classical interdigital filters based on cylindrical posts is also discussed, with the aim of  
 107 comparing both topologies in terms of compactness, electrical performance and power  
 108 handling capability.

## 109 2 Design of the coaxial resonator and related equivalent circuit

110 The basic coaxial resonator proposed in this work is shown in Fig. 2. It is based  
 111 on a rectangular cavity loaded with a thick metallic strip. The strip is connected to the  
 112 base of the resonator at one end. At the other end, there is an open circuit or gap, as  
 113 shown in Fig. 2. In addition, the metal strip is located in the center of the rectangular  
 114 enclosure. The relevant dimensions of the resonator are designed with the objective of  
 115 having a resonance at the center frequency  $f_0$  of the filter. In this work, two bandpass



**Figure 1.** Third-order symmetric interdigital bandpass filter with flat metallic strips.

116 filters of orders  $N=3$  and  $N=5$  will be designed. Since both of them share the same  
 117 center frequency ( $f_0 = 3$  GHz), the resonator designed in this section will be used in  
 118 both bandpass filters.

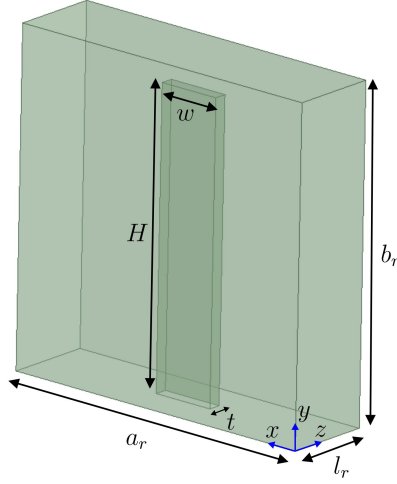
119 In order to design the resonator, the full-wave EM simulator Ansys HFSS has been  
 120 used. First of all, an access port located at the  $y = 0$  plane of the structure in Fig. 2,  
 121 has been opened. It is important to note that, with this access port, the resonator be-  
 122 comes a section of coaxial guide, with internal and external rectangular contours, loaded  
 123 by a length of short-circuited rectangular waveguide. To continue, the phase of the  $S_{11}$   
 124 parameter at the input port is calculated, and the dimensions of the resonator are opti-  
 125 mized until the phase of  $S_{11}$  is equal to  $180^\circ$  at a frequency equal to  $f_0$ . The final di-  
 126 mensions (all expressed in mm) of the optimized resonator are as follows (see Fig. 2):  
 127  $a_r = b_r = 26$ ,  $l_r = 7$ ,  $H = 23.233$ ,  $w = 5$  and  $t = 1$ . It is important to note that the di-  
 128 mensions  $a_r, b_r, l_r$  (dimensions of the rectangular enclosure),  $w$  and  $t$  (width and thick-  
 129 ness of the strip) will remain fixed, and will not be optimized during the filter design pro-  
 130 cess.

131 In addition, we are also interested in obtaining an equivalent circuit of the resonator  
 132 based on lumped elements. This equivalent circuit consists of a series combination of an  
 133 inductance  $L$  and a capacitor  $C$ . With the aim of obtaining the equivalent series net-  
 134 work, the slope parameter  $X$  of the optimized coaxial resonator is first calculated from  
 135 the simulated results (Matthaei et al., 1980):

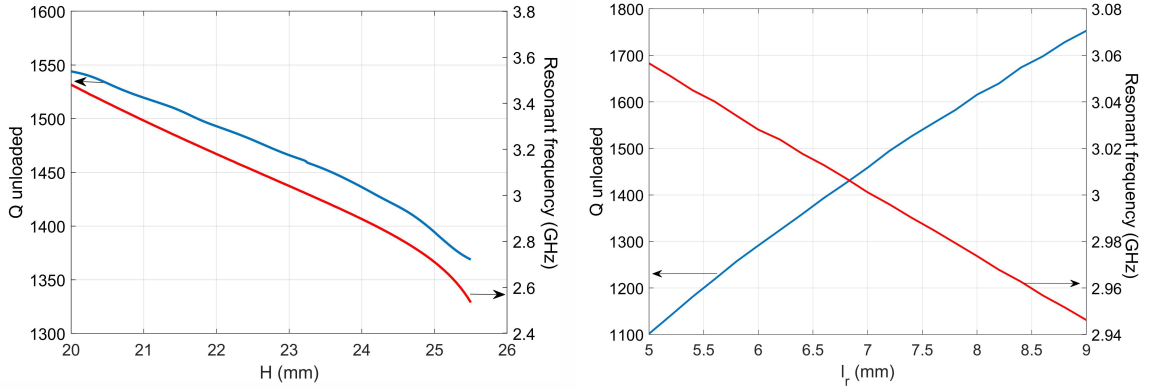
$$X = \frac{w_0}{2} \left. \frac{dX_{in}(w)}{dw} \right|_{w=w_0} = 295.48 \Omega \quad (1)$$

136 where  $w_0 = 2\pi f_0$ , and  $X_{in}$  represents the reactance of the input impedance of the op-  
 137 timized resonator. Then, the value of the equivalent inductance of the resonator can be  
 138 calculated as:  $L = X/w_0 = 15.676$  nH. The value of the equivalent capacitance can  
 139 now be computed using:  $C = 1/(w_0^2 L) = 0.179$  pF.

140 The unloaded  $Q$  factor of the designed resonator has been calculated using Ansys  
 141 HFSS. In this analysis, an equivalent conductivity value of  $38 \cdot 10^6$  S/m (aluminium)  
 142 has been considered, thus obtaining  $Q_u = 1462$ . Furthermore, we have computed both  
 143 the unloaded  $Q$  factor and the resonant frequency ( $f_0$ ) of the designed resonator in terms  
 144 of the height  $H$  of the strip, and also as a function of the length  $l_r$  of the rectangular en-  
 145 closure. The obtained results are shown in Fig. 3. As we expected, when the height of  
 146 the strip decreases (thus increasing the gap between the strip and the top wall of the en-  
 147 closure), both  $Q_u$  and  $f_0$  increase (see Fig. 3, left). As it is well known, the cited gap  
 148 represents a critical region in terms of high-power performance in classical coaxial res-  
 149 onator filters. Nevertheless, since the thickness of the strip used in our proposed resonator



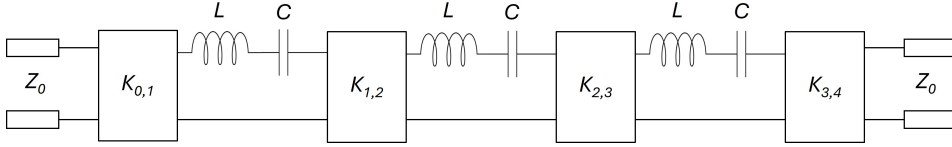
**Figure 2.** Coaxial resonator loaded with a flat metallic strip.



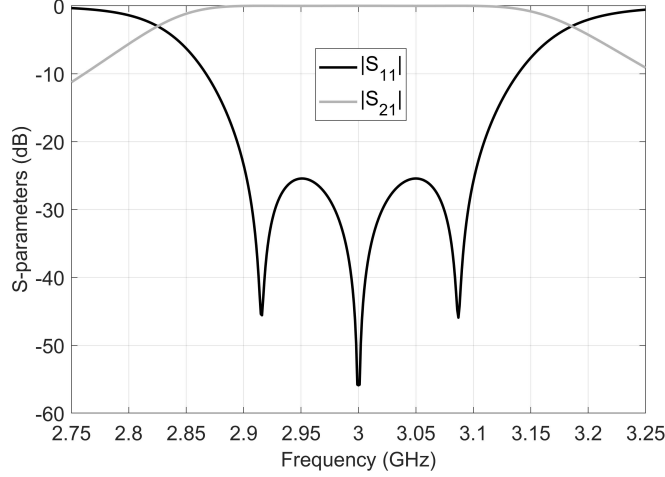
**Figure 3.** Variation of the unloaded  $Q$  factor and the resonant frequency of the flat strip resonator in terms of  $H$  and  $l_r$ .

150 is small (only 1 mm), the equivalent capacitance in this region is less significant than the  
 151 one found in classical resonators based on cylindrical posts. Therefore, an increased high-  
 152 power performance is expected compared to classical interdigital filters, as it is discussed  
 153 in Section 6.

154 On the other hand, as it can be observed in Fig. 3 (right), when the length of the  
 155 rectangular enclosure decreases (thus obtaining a more compact resonator),  $Q_u$  decreases  
 156 (thus increasing the losses, as well), while  $f_0$  increases. Furthermore, for comparison pur-  
 157 poses, we have also analysed the unloaded  $Q$  factor of a classical resonator loaded with  
 158 a cylindrical post. In this design, the post diameter has been chosen to be equal to the  
 159 width  $w$  of the flat strip, while the dimensions of the rectangular enclosure are the same  
 160 that we have used in the flat strip resonator. The computed  $Q$  factor is  $Q_u = 937$ , so  
 161 the novel resonator proposed in our work has a better performance in terms of losses than  
 162 the classical solution based on cylindrical posts.



**Figure 4.** Equivalent circuit network of the 3rd-order interdigital bandpass filter.



**Figure 5.** Frequency response of the equivalent circuit network shown in Fig. 4.

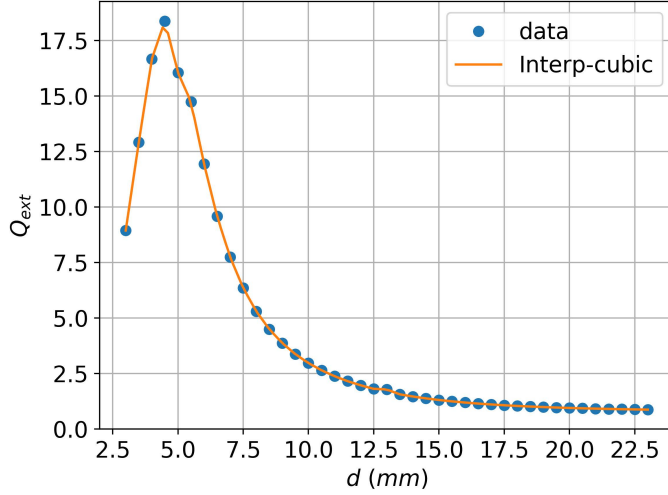
### 3 Equivalent circuit of the bandpass filters

The purpose of this section is to obtain an equivalent network to characterize the designed bandpass filters. The electrical response of this equivalent circuit will be used to generate target curves to obtain the in-band response of the designed filter using iterative optimizations. The equivalent circuit of the filters will be composed of the cascade connection of ideal impedance inverters (which will characterize both the input/output couplings and also the inter-resonator couplings) and series L-C networks (which will model the resonators). It is also worth noting that, as it is well known, transmission zeros cannot be realized using coaxial resonators in a standard in-line interdigital configuration.

The electrical specifications for the 3rd-order interdigital filter that will be designed in Section 4 are: center frequency  $f_0 = 3$  GHz, return loss of  $RL = 25$  dB, and bandwidth  $BW = 200$  MHz. Given these specifications, the values (in  $\Omega$ ) of the impedance inverters, which have been calculated using the slope parameter  $X$  of the resonator, are as follows:  $K_{01} = K_{34} = 86.394$  and  $K_{12} = K_{23} = 24.027$  (Matthaei et al., 1980). The equivalent circuit of the 3rd-order filter is shown in Fig. 4, while its frequency response is displayed in Fig. 5.

### 4 Design of a 3rd-order bandpass filter following a classical approach

In this section, the procedure used to design a third-order Chebyshev bandpass filter (electrical specifications:  $f_0 = 3$  GHz,  $RL = 25$  dB,  $BW = 200$  MHz) is described in detail. As we have previously mentioned, in this case, where the order of the filter is not high, the design process is based on the classical approach proposed in (Cameron et al., 2018). First, input and output couplings are characterized and, after that, the inter-resonator couplings are calculated. Finally, the complete filter is assembled and all di-



**Figure 6.** External quality factor in terms of the distance  $d$ .

186 mensions are optimized using as a target the response of the corresponding equivalent  
 187 circuit (see Fig. 5).

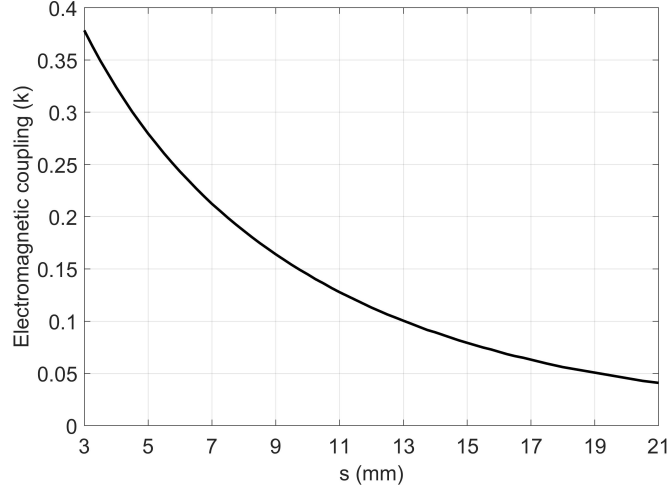
188 First, the design of the excitation of the input and output resonators is discussed.  
 189 These resonators are fed with SMA coaxial waveguides ( $\epsilon_r = 2.2$ ), with external and in-  
 190 ternal radii of 2.05 and 0.65 mm, respectively. The metallic strips of the input/output  
 191 resonators are in contact with a metal feeding probe, located at a distance  $d$  from the  
 192 bottom of the cavity, as shown in Fig. 1. The distance  $d$  will be optimized in order to  
 193 obtain the required external quality factor  $Q_{ext}$ . This parameter can be obtained follow-  
 194 ing the method detailed in (Cameron et al., 2018), namely, by evaluating the group de-  
 195 lay of parameter  $S_{11}$  at frequency  $f_0$ . Fig. 6 shows the dependence of the external qual-  
 196 ity factor in terms of the distance  $d$ . On the other hand, the required value for  $Q_{ext}$  can  
 197 be calculated as follows (Cameron et al., 2018):

$$Q_{ext} = \frac{f_0}{BW m_{01}^2} = 10.055 \quad (2)$$

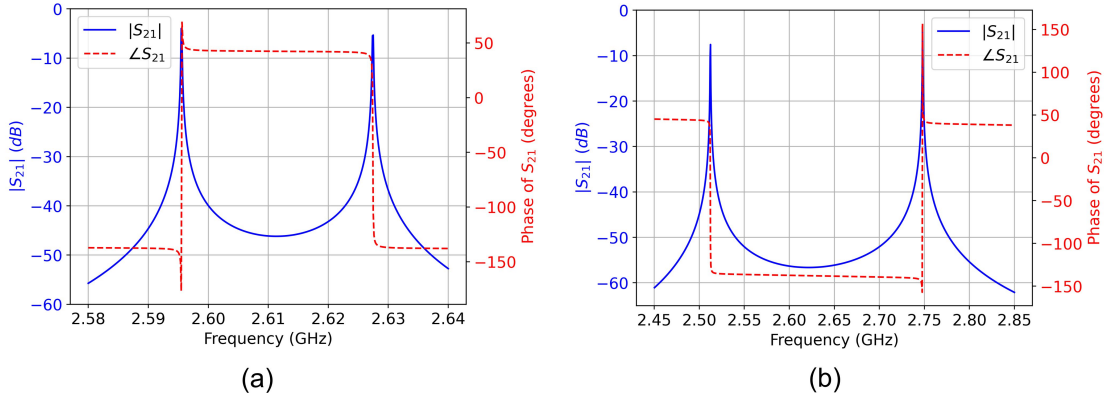
198 where  $m_{01} = 1/\sqrt{g_0 g_1}$  (being  $g_0$  and  $g_1$  the corresponding low-pass prototype elements).  
 199 Furthermore,  $Q_{ext}$  is related to  $K_{01}$  through  $Q_{ext} = (X Z_0)/K_{01}^2$ , where  $Z_0 = \eta_0/\sqrt{\epsilon_r} =$   
 200  $253.9 \Omega$  represents the modal impedance of the TEM mode of the coaxial waveguide port.  
 201 In view of the value obtained for the required external quality factor, an initial value  $d =$   
 202  $7.321$  mm can be deduced from Fig. 6.

203 Once the couplings of the input/output resonators have been obtained, the cou-  
 204 plings between adjacent resonators are characterized as a function of the spacing  $s$  be-  
 205 tween the strips (see Fig. 1). For this purpose, a symmetrical two-port component is con-  
 206 sidered, consisting of two resonators weakly coupled to the input and output coaxial ports.  
 207 It is well known that the resonant frequencies of this component are related to the value  
 208 of the electromagnetic coupling  $k$  of the structure, which can be obtained as a function  
 209 of the separation  $s$  between strips, as shown in Fig. 7. Since the required value for the  
 210 electromagnetic coupling between resonators #1 and #2 can be computed as  $k_{12} = K_{12}/(w_0 L) =$   
 211  $0.0813$ , an initial value of  $s = 14.774$  mm can be deduced from Fig. 7.

212 At this point, it would be interesting to determine the nature of the coupling (elec-  
 213 tric or magnetic) between two interdigital resonators based on flat metallic strips. First,



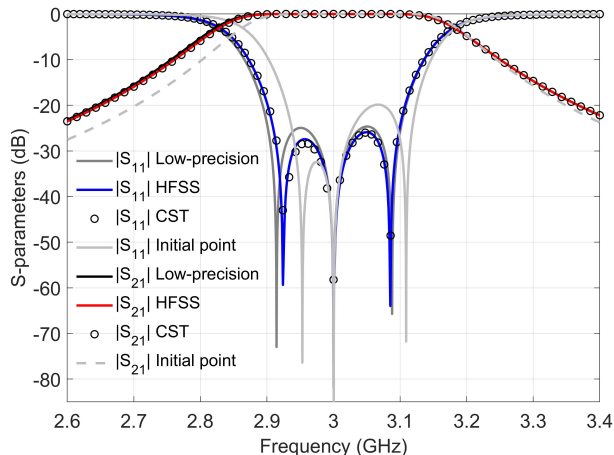
**Figure 7.** Electromagnetic coupling between adjacent resonators as a function of the separation  $s$  between strips.



**Figure 8.** Magnitude and phase of the  $S_{21}$  parameter when the resonators are in (a) a combline configuration and in (b) an interdigital configuration.

214 let us consider two resonators coupled in a combline configuration, both weakly coupled  
 215 to the coaxial input/output ports (the dimensions of the strips and the dimensions of  
 216 the resonator have been deduced in Section 2). In this component, note that both res-  
 217 onator strips are connected to the bottom of the filter. In addition, a separation of  $s =$   
 218  $14$  mm between the resonators has been used for this simulation. It is well-known that,  
 219 in this type of structure, the nature of the predominant coupling is magnetic (Thomas,  
 220 2003). Moreover, Fig. 8(a), which can be used to determine the magnitude of the elec-  
 221 tromagnetic coupling  $k$  (Cameron et al., 2018), shows the magnitude and the phase of  
 222 the  $S_{21}$  parameter of the considered combline structure. On the other hand, Fig. 8(b)  
 223 shows this same information when the two resonators are in an interdigital configura-  
 224 tion. As we can observe, there is a phase shift of  $180^\circ$  between the  $S_{21}$  phase responses  
 225 of the two cases analyzed, thus confirming a change in the sign of the predominant cou-  
 226 pling. In this regard, it is important to note that the combline-type coupling (of mag-  
 227 netic nature) is approximately one order of magnitude lower than the interdigital one  
 228 for the same resonator separation, thus making the non-adjacent couplings negligible in  
 229 the interdigital case and, thereby, facilitating the proposed systematic in-line design.



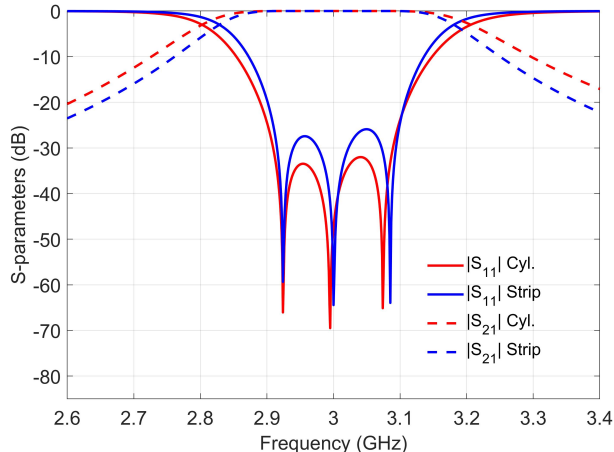


**Figure 9.** S-parameters of the designed 3rd-order interdigital filter. Comparison of results obtained with Ansys HFSS and CST Studio Suite.

230 In the next step of the design process, the complete filter is assembled (considering  
 231 symmetry in the component) and the filter dimensions  $H_1$ ,  $H_2$ ,  $s$  and  $d$  (see Fig. 1)  
 232 are optimized in a low-accuracy space, until recovering the target response. The values  
 233 obtained are (all in mm)  $H_1 = 23.380$ ,  $H_2 = 20.144$ ,  $s = 18.606$  and  $d = 5.491$ . All sim-  
 234 ulations previously described have been performed using the full-wave EM tool FEST3D  
 235 (which is part of CST Studio Suite), with a set of parameters chosen to obtain very fast  
 236 low-accuracy results. For instance, only 10 accessible modes (Conciauro et al., 2000) were  
 237 considered in all simulations concerning the design process of this filter. Consequently,  
 238 the resulting low-accuracy frequency response is not sufficiently precise for fabrication  
 239 purposes. However, the well-known ASM technique presented in (Bandler et al., 1995)  
 240 allows us to refine these results by transforming the low-accuracy data (coarse model)  
 241 into a high-precision solution (fine model). In this design, FEST3D serves as the coarse  
 242 model, while the finite-element EM tool Ansys HFSS is used as the fine model. In most  
 243 cases, only a few iterations are required to obtain the target frequency response of the  
 244 filter in the fine space.

245 The final dimensions obtained with the ASM procedure (in mm) are as follows:  $H_1 =$   
 246  $23.683$ ,  $H_2 = 20.320$ ,  $s = 17.721$  and  $d = 6.076$ . In order to further validate both the  
 247 topology of the new filter and the design procedure we have followed, the filter has been  
 248 analyzed using an additional EM simulator, namely, CST Studio Suite. The results ob-  
 249 tained are shown in Fig. 9, where an excellent agreement between the results provided  
 250 by HFSS and CST can, indeed, be observed. We have also included in the same figure  
 251 the S-parameters of the filter designed in the low-precision space to show that a very good  
 252 agreement between the results obtained in both simulation spaces (low and high-precision  
 253 models) has been achieved. Furthermore, the response in the fine space with the opti-  
 254 mum dimensions found in the coarse space has also been included in Fig. 9 for compar-  
 255 ison purposes (see label “Initial point”).

256 At this point, it is important to discuss the design of classical interdigital filters based  
 257 on cylindrical posts, with the aim of comparing this classical topology to the novel pro-  
 258 posed flat strip configuration, both in terms of compactness and electrical performance.  
 259 In this context, we have first designed a resonator ( $f_0 = 3$  GHz) loaded with a cylin-  
 260 drical post whose diameter is equal to the width  $w$  of the flat strip that we have used  
 261 in our previous design (the dimensions of the rectangular enclosure used in this resonator  
 262 can be found in Section 2). Afterwards, the inter-resonator couplings have been com-



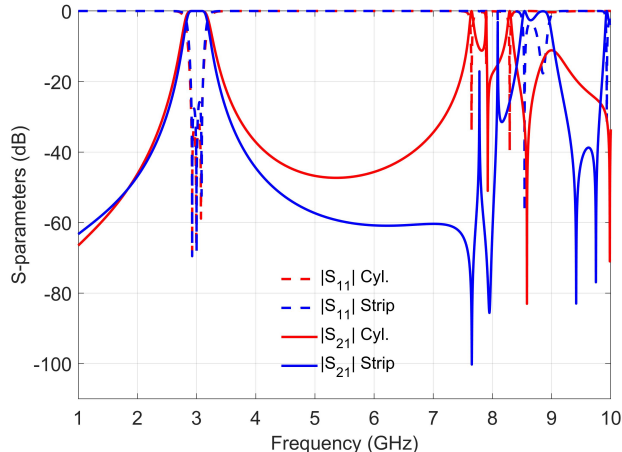
**Figure 10.** Electrical responses of the designed 3rd-order interdigital bandpass filters based on cylindrical posts and flat strips.

263 puted to determine the initial values for the different separations between resonators.  
 264 Our results show that such initial values are very similar to the ones obtained when res-  
 265 onators based on flat strips are considered. We can, therefore, estimate that the final fil-  
 266 ter lengths for both considered topologies are comparable, so the compactness of our novel  
 267 filter topology based on flat strips resonators is very similar to the one obtained with clas-  
 268 sical resonators loaded with cylindrical posts.

269 To validate this statement, a 3rd-order interdigital bandpass filter including res-  
 270 onators based on cylindrical posts has been designed following the approach previously  
 271 detailed. The only difference is that, in this case, after computing the initial values for  
 272 all relevant filter dimensions, a brute-force optimization algorithm has been used (instead  
 273 of the ASM technique) in order to obtain the filter dimensions in a high-precision space.  
 274 In this design, we have used the same coaxial excitation and the same electrical spec-  
 275 ifications that we employed for the 3rd-order filter based on flat strips. The dimensions  
 276 of this filter (all in mm) are:  $H_1 = 22.70$ ,  $H_2 = 19.56$ ,  $s = 18.54$  and  $d = 5.40$ , where  $s$   
 277 represents the separation between the centres of the cylindrical posts (see also Fig. 1 for  
 278 the definition of these variables by considering cylindrical posts). The total length of the  
 279 designed filter (high-precision space) is 44.1 mm, while the length of the 3rd-order fil-  
 280 ter based on flat strips previously designed is equal to 44.4 mm, thus confirming that both  
 281 topologies exhibit an almost identical performance in terms of compactness. The in-band  
 282 frequency responses (obtained using Ansys HFSS) of both 3rd-order filter topologies are  
 283 shown in Fig. 10.

284 Regarding the insertion loss, we have considered a finite conductivity of  $38 \cdot 10^6$   
 285 S/m (aluminium) in our analysis, and we have obtained that the insertion loss mean value  
 286 is equal to 0.1 dB in the nominal passband (2.9 – 3.1 GHz) for both topologies. It is im-  
 287 portant to note at this point that, although the two filters have been designed with the  
 288 same electrical specifications, the bandwidth of the filter based on cylindrical posts gets  
 289 slightly higher than the bandwidth of the filter based on flat strips (see also Fig. 10). This  
 290 fact explains why both filters present a very similar insertion loss, even though the res-  
 291 onator based on flat strips exhibits a higher unloaded  $Q$  factor.

292 We have also calculated the insertion loss of the 5th-order interdigital filter based  
 293 on flat strips presented in the next section. In this case, the mean value in the nominal  
 294 passband is of 0.24 dB. It is worth noting that the obtained values are consistent with



**Figure 11.** Broadband response of the designed 3rd-order interdigital bandpass filters.

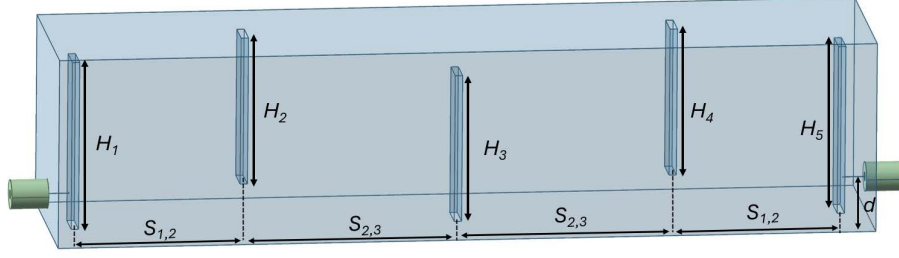
295 the theoretical insertion loss computed using the analytical formulas found in the tech-  
 296 nical literature for the prescribed specifications.

297 On the other hand, and regarding the broadband performance, Fig. 11 shows the  
 298 out-of-band response of the designed 3rd-order interdigital filter based on cylindrical posts  
 299 (red lines) compared to the response of the filter based on flat strips (blue lines). It is  
 300 worth mentioning that the interdigital filter based on flat strips shows an out-of-band  
 301 rejection starting at higher frequencies, and a wider spurious-free range, when compared  
 302 to the classical filter based on cylindrical posts. Therefore, the novel topology proposed  
 303 in this work enhances the out-of-band performance of classical interdigital filters.

304 As we have mentioned before, a 3rd-order interdigital bandpass filter including res-  
 305 onators based on cylindrical posts has been designed using a brute-force algorithm to  
 306 obtain the filter dimensions in a high-precision space. In particular, after obtaining an  
 307 initial set of filter dimensions following the classical design approach, the whole compo-  
 308 nent is assembled. Next, by means of a full-wave simulator (Ansys HFSS), the dimen-  
 309 sions of the physical structure are iteratively optimized until the desired electrical re-  
 310 sponse is recovered. Proceeding in this way, the complete set of the filter dimensions (4  
 311 dimensions in this case) is optimized at the same time. Compared to the design strat-  
 312 egy based on the ASM technique, the design approach based on a brute-force design al-  
 313 gorithm spends, approximately, 19% more computational time. In addition, it is worth  
 314 noting that this conclusion stands for a third-order interdigital filter. Therefore, if a high-  
 315 order filter needs to be designed, the number of involved physical dimensions will be higher  
 316 and the cited increase in CPU time could be even more significant, since the brute-force  
 317 algorithm does not benefit from the segmentation technique proposed in this work.

## 318 5 Design of a 5th-order bandpass filter using a segmentation technique

319 In this section, the design of a fifth-order interdigital bandpass filter with a Cheby-  
 320 shev response is addressed. The technique proposed for the efficient design of this high-  
 321 order filter is based on a segmentation of the component, where the different resonators  
 322 of the filter are progressively added one after the other, with the aim of reducing the num-  
 323 ber of variables to be optimized in each step of the design process. The electrical response  
 324 (i.e., the group delay of the  $S_{11}$  parameter) of the equivalent circuit of the structure used  
 325 in each step, will be used as the target curve in the optimization process. After that, the



**Figure 12.** Fifth-order interdigital bandpass filter with flat metallic strips.

326 ASM technique will be used again to obtain the final dimensions. Fig. 12 shows the topol-  
 327 ogy of the 5th-order filter under consideration.

328 The center frequency chosen for this bandpass filter is the same as the one we used  
 329 for the 3rd-order filter. Therefore, both the resonator dimensions and the values of the  
 330 equivalent L-C series network have been already calculated in Section 2. The equivalent  
 331 circuit of this filter, as explained in Section 3, consists of the cascade connection of L-  
 332 C series networks and ideal impedance inverters. Given the electrical specifications of  
 333 the 5th-order filter ( $f_0 = 3$  GHz,  $RL = 25$  dB,  $BW = 150$  MHz), the following values  
 334 (in  $\Omega$ ) of the impedance inverters can be readily deduced:  $K_{01} = K_{56} = 68.658$ ,  $K_{12} =$   
 335  $K_{45} = 14.387$ ,  $K_{23} = K_{34} = 10.083$ .

### 336 5.1 Design of the coaxial excitation

337 First, the design of the coaxial excitation of the input/output resonators of the band-  
 338 pass filter is addressed. The dimensions of the coaxial waveguide ( $\epsilon_r = 2.2$ ) are the same  
 339 ones used in the 3rd-order filter (external and internal radii of 2.05 and 0.65 mm, respec-  
 340 tively). In order to characterize the external quality factor  $Q_{ext}$ , we follow the same method  
 341 explained for the 3rd-order filter. In this case, since the required quality factor is  $Q_{ext} =$   
 342  $15.921$ , an initial value of  $d = 5.053$  mm can be deduced from Fig. 6.

### 343 5.2 Inter-resonator coupling

344 The inter-resonator couplings depend on the separation  $s_{i,j}$  between the resonator  
 345 strips (see Fig. 12). The required electromagnetic couplings between resonators can be  
 346 computed as:

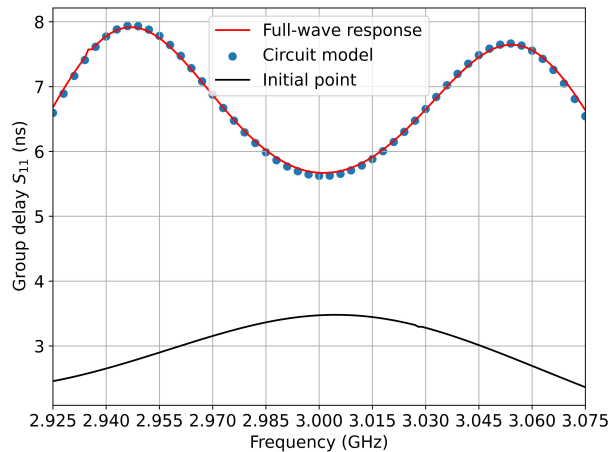
$$k_{1,2} = \frac{K_{1,2}}{w_0 \cdot L} = 0.0487 \quad (3)$$

$$k_{2,3} = \frac{K_{2,3}}{w_0 \cdot L} = 0.0341 \quad (4)$$

347 Therefore, the following initial values for the separation between resonators can be de-  
 348 duced from Fig. 7:  $s_{1,2} = 19.236$  mm and  $s_{2,3} = 22.551$  mm.

### 349 5.3 Systematic design using a step-by-step procedure

350 Once we have obtained a set of initial values for the relevant dimensions, we con-  
 351 tinue our filter design following a step-by-step strategy. The design process is based on  
 352 the segmentation of the filter, where the different resonators are added one after the other,



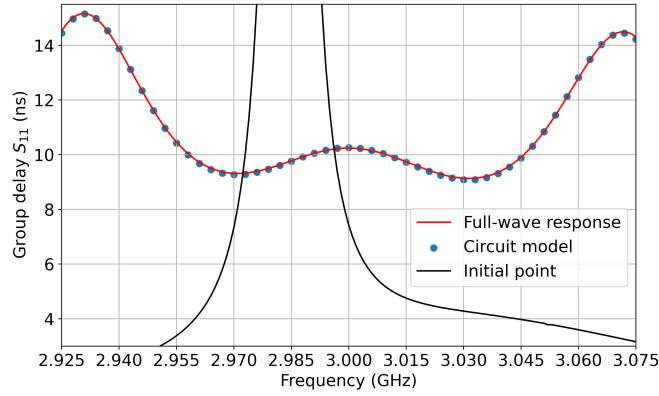
**Figure 13.** Electrical response of the waveguide structure considered in the first stage compared to the ideal response.

353 with the aim of reducing the number of variables to be optimized in each stage. The elec-  
 354 trical response of the lumped-element equivalent circuit of the structure used in each step  
 355 will be used as the target response. The optimized dimensions obtained in each step will  
 356 then be used as initial values in the next step of the design process. This design proce-  
 357 dure will be concluded with the ASM technique. All the simulations discussed in the next  
 358 sections have been performed in a low-precision space using the EM software FEST3D.

### 359 5.3.1 First stage

360 In the first step of the design process, the component under consideration consists  
 361 of the first two resonators of the interdigital filter (see Fig. 12) fed by a coaxial probe.  
 362 It is worth noting that there is only one access port in this structure. The equivalent cir-  
 363 cuit of this waveguide component consists of the cascade connection of the input trans-  
 364 mission line, inverter  $K_{0,1}$ , an L-C series network representing the first resonator, impedance  
 365 inverter  $K_{1,2}$ , and another L-C series network short-circuited at its end. In this first stage,  
 366 the dimensions to be optimized are the heights  $H_1$  and  $H_2$  of the strips, the separation  
 367  $s_{1,2}$  between resonators, and the height  $d$  of the feeding probe (see Fig. 12). In this first  
 368 stage, the set of dimensions obtained in the previous sections are used as a starting point  
 369 in the optimization process.

370 Next, the dimensions of the waveguide component are optimized using, as a tar-  
 371 get response, the group delay of the  $S_{11}$  parameter of the circuit model. The results of  
 372 the optimization are shown in Fig. 13, where a very good agreement can be observed be-  
 373 tween both set of data. The values obtained (in mm) for each dimension at the end of  
 374 the optimization process are:  $H_1 = 22.942$ ,  $H_2 = 22.205$ ,  $s_{1,2} = 19.059$  and  $d = 4.522$ .  
 375 On the other hand, the set of initial dimensions used as a starting point in this first stage  
 376 are:  $H_1 = H_2 = 23.233$ ,  $s_{1,2} = 19.236$  and  $d = 5.053$  (the response of this initial point  
 377 has also been included in Fig 13). As we can see, the optimized final values are, indeed,  
 378 very close to the initial values.



**Figure 14.** Electrical response of the waveguide structure considered in the second stage compared to the ideal response.

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### 5.3.2 Second stage

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In the second stage of the design process, a new resonator is added to the structure, thus obtaining a waveguide structure composed of the first three resonators of the filter (see Fig. 12). The circuit model of this structure can be readily obtained by adding to the circuit network of the previous step, the inverter  $K_{2,3}$ , and a new L-C series network. Next, the dimensions of this waveguide structure must be optimized to obtain the same electrical response of the ideal network. The parameters to be optimized in this stage are the same as in the previous step ( $H_1$ ,  $H_2$ ,  $s_{1,2}$  and  $d$ ), including also the height  $H_3$  of the third resonator, and the separation  $s_{2,3}$  between the second and third resonators (see also Fig. 12). The results of this optimization are shown in Fig. 14, where an excellent agreement is again observed between the response of the lumped-element model and the full-wave simulation. The final values (in mm) for the optimized dimensions are:  $H_1 = 22.955$ ,  $H_2 = 20.108$ ,  $H_3 = 21.994$ ,  $s_{1,2} = 22.226$ ,  $s_{2,3} = 25.176$  and  $d = 4.382$ . The corresponding starting points are the final values obtained at the end of the previous stage. We have also included in Fig. 14 the response of the initial point from where the matching to the group delay reference curve is performed.

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### 5.3.3 Third stage

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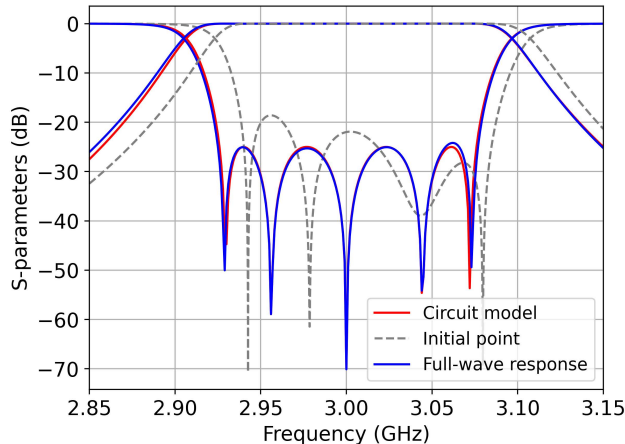
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At this point of the procedure, we have completed the design of one half of the band-pass filter, since the component is symmetrical. Therefore, the optimization of the whole filter can be performed next by taking into account that (see Fig. 12):  $H_1 = H_5$ ,  $H_2 = H_4$ ,  $s_{1,2} = s_{4,5}$  and  $s_{2,3} = s_{3,4}$ . In this final stage of the design process, the S-parameters of the equivalent lumped-element network of the 5th-order filter are used as the target response, and all the dimensions of the structure are optimized. The optimization results are shown in Fig. 15, where an excellent agreement is observed between the full-wave EM-simulation, and the response of the ideal network, thereby validating the systematic design procedure proposed in this work. The final dimensions (in mm) of the filter (obtained in a low-precision space) are:  $H_1 = H_5 = 22.960$ ,  $H_2 = H_4 = 20.099$ ,  $H_3 = 20.069$ ,  $s_{1,2} = s_{4,5} = 22.198$ ,  $s_{2,3} = s_{3,4} = 28.463$  and  $d = 4.382$ .



**Figure 15.** S-parameters of the 5th-order interdigital filter (coarse model) and response of the coarse model filter in the high-precision space (label: initial point).

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#### 5.4 Implementation of the ASM Technique

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As already mentioned, all previous simulations have been performed in a low-accuracy space (using FEST3D). To continue, therefore, the final dimensions obtained in the previous section need to be mapped to a high-precision space using the ASM method. In this context, we have included in Fig. 15 the filter response in the high-precision space using the optimum dimensions designed in the coarse space.

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Fig. 16 shows the frequency response of the filter after implementing the ASM technique. Again, two different EM simulators (Ansys HFSS and CST Studio Suite) have been used to validate the design process, obtaining an excellent agreement between both set of simulated data. In addition, the electrical response of the filter designed in the low-accuracy simulation space has also been included in Fig. 16. The values (in mm) of the dimensions of the 5th-order interdigital filter (after implementing the ASM technique) are:  $H_1 = H_5 = 23.100$ ,  $H_2 = H_4 = 20.204$ ,  $H_3 = 20.174$ ,  $s_{1,2} = s_{4,5} = 21.839$ ,  $s_{2,3} = s_{3,4} = 28.059$  and  $d = 4.421$ .

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Finally, we have compared in Table 1 the electrical performance of the bandpass filters designed in this work (3rd and 5th-order filters) with other similar designs (found in the technical literature) concerning interdigital waveguide bandpass filters using resonator posts with a cylindrical or a rectangular geometry.

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## 6 Multipactor breakdown prediction

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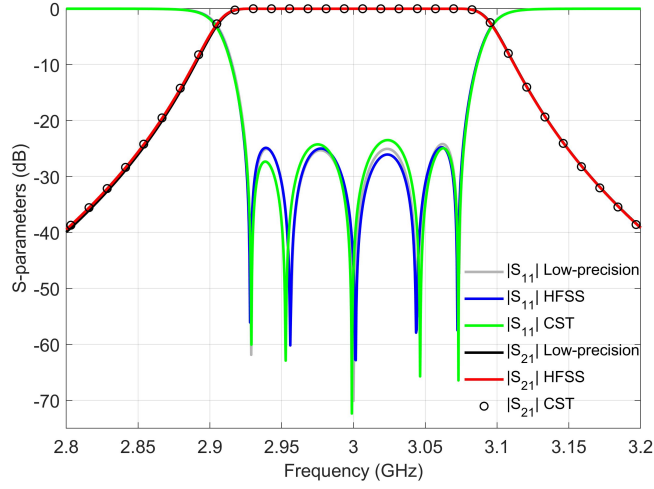
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A power handling capability study of the designed interdigital filters is addressed in this section. In particular, the high-power performance of both topologies of interdigital filters discussed in this work (based on cylindrical posts and on flat strips) will be compared. The multipactor effect is a high-power phenomenon that appears in components used in communication satellites, under high vacuum conditions and in the presence of high-intensity electromagnetic fields, linked to an electron avalanche effect (Vaughan, 1988). This phenomenon can lead to a degradation of the device performance, so it is very important to properly characterize the behaviour of the designed components under high-power conditions.

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In this context, Fig. 17 shows the electric field distribution in the 5th-order interdigital filter based on flat strips computed at  $f_0 = 3$  GHz (with values scaled for an in-



**Figure 16.** S-parameters of the 5th-order interdigital bandpass filter in the high-precision space. Comparison of results obtained with Ansys HFSS and CST Studio Suite.

**Table 1.** Electrical performance of several interdigital waveguide bandpass filters found in the technical literature.

Ref.	$f_0$ (GHz)	$N$	$BW$ (GHz)	$RL$ (dB)	Post geometry
<b>This work</b>	3	3	0.2	25	Rect.
<b>This work</b>	3	5	0.2	25	Rect.
(Packiaraj et al., 2005)	0.975 - 1.825	3	0.25 / 0.35	15	Cyl.
(Ikram & Rana, 2006)	6.175	5	0.5	10	Cyl.
(Munina & Turalchuk, 2015)	2.91 - 3 - 3.09	3	$\simeq 0.04 - 0.04 - 0.04$	20	Cyl.
(J. Li et al., 2018)	15	8	4.4	10	Cyl.
(Wu et al., 2019)	14	4	1.4	15	Rect.
(Jiang et al., 2019)	11.725	8	2.05	17	Rect.
(Rodriguez-Morales et al., 2021)	3.75	11	2	10	Rect.
(Rodriguez-Morales et al., 2021)	6	11	2	10	Rect.
(Kong et al., 2023)	3.8	4	0.9	18	Rect.



put power level of 1 W). The electric field distribution reveals a high concentration of the electric field around all the metallic strips. Typically, the critical areas in terms of multipactor effect for classical interdigital resonators based on cylindrical posts are the zones between the top surface of the post and the top/bottom walls of the rectangular enclosure. Nevertheless, for the novel resonators based on planar strips proposed in this work, since the thickness of the strip used in our designs is small (only 1 mm), the equivalent capacitance in this region is less significant than the one found in classical resonators based on cylindrical posts. Therefore, a multipactor analysis in several isolated regions of the designed filters is required to identify the critical zones in this new filter topology.

After analysing the evolution of the electrons trajectory in the designed filters, we have identified that the critical areas in these structures are the flat surfaces of the first and last resonator strips facing the front and back walls of the rectangular enclosure (where the distance is of 3 mm). In these zones, there is, indeed, a large capacitance (probably the largest in the whole structure), thus becoming the most critical regions in terms of the multipactor phenomenon. This fact has to be taken into account if these filters are to be used in space applications, since, under vacuum conditions, they will be susceptible to resonant electron discharges.

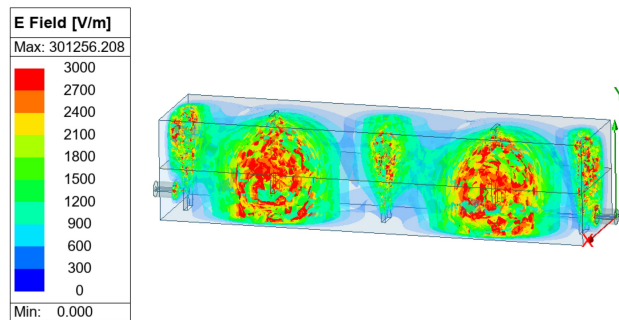
Next, a power analysis of the interdigital filters designed in this work (considering that they have been manufactured in aluminium) has been performed under vacuum conditions using the commercial software Spark3D (which is part of CST Studio Suite). To do this, we have first used Ansys HFSS to compute the EM fields at the central frequency of the designed filters ( $f_0 = 3$  GHz in all cases), and also at the frequencies ( $f_1$  and  $f_2$ ) where the  $S_{21}$  parameter group delay and stored energy are maximum. Next, the calculated EM fields are imported by the Spark3D simulator, which is used to estimate the multipactor power threshold ( $P_{th}$ ) of the designed filters. A homogeneous electron seeding of 10000 initial electrons has been defined in the simulations. In addition, the main parameters of the SEY curve are:  $\delta_{\max} = 2.92$ ,  $\delta_0 = 0.8$ ,  $E_1 = 17$  eV and  $E_{max} = 276$  eV (ESA-ESTEC, E. P. D., 2020).

The obtained results for the 3rd-order interdigital filter based on flat strips are ( $f_1 = 2.86$  GHz,  $f_2 = 3.15$  GHz):  $P_{th,f_0} = 170$  W,  $P_{th,f_1} = 72.8$  W and  $P_{th,f_2} = 83.5$  W. Moreover, for the 5th-order interdigital filter we have obtained ( $f_1 = 2.91$  GHz,  $f_2 = 3.09$  GHz):  $P_{th,f_0} = 100$  W,  $P_{th,f_1} = 45.4$  W and  $P_{th,f_2} = 51.3$  W. Although the obtained multipactor power thresholds are lower at frequencies  $f_1$  and  $f_2$ , in practice, the frequency bandwidth of the employed signals do not usually reach such extreme values. The minimum of these computed  $P_{th}$  values would be the maximum operating power of each filter in space applications. If needed, solutions like the one proposed in (Coves et al., 2023) can be used to increase these values.

In order to compare the power handling capability of the new proposed topology with that of the classical solution based on cylindrical posts, a multipactor analysis has also been performed for the designed 3rd-order interdigital filter based on cylindrical resonators. The obtained results are ( $f_1 = 2.83$  GHz,  $f_2 = 3.17$  GHz):  $P_{th,f_0} = 79.4$  W,  $P_{th,f_1} = 40.9$  W and  $P_{th,f_2} = 31.4$  W. Therefore, compared to the data obtained for the 3rd-order filter based on flat strips, we can conclude that the new proposed topology exhibits a better high-power performance than the classical interdigital filters based on cylindrical resonators.

## 7 Conclusion

A systematic procedure for the efficient design of a new topology for interdigital bandpass filters, based on flat metallic strips, has been proposed. The design method combines the use of the classical circuit-based approach and a strategy based on the seg-



**Figure 17.** Electric field ( $f_0 = 3$  GHz) in the 5th-order interdigital bandpass filter.

487 mentation of the filter, with the aim of reducing the number of variables to be optimized  
 488 at each stage of the design process. Moreover, the ASM technique has been used in order  
 489 to perform most of the simulations in a low-accuracy space, thus reducing the overall  
 490 computational effort. Two interdigital bandpass filters operating in the S-band have  
 491 been designed, and their electrical responses have been successfully compared to the results  
 492 provided by two different EM full-wave simulators, thus validating both the new  
 493 topology and the design procedure proposed. In addition, a multipactor study has been  
 494 performed, in order to characterize the power capability of the components under vacuum  
 495 conditions, in view of a possible use of the designed filters in satellite payloads. Finally,  
 496 the design of a 3rd-order interdigital filter based on cylindrical posts has also been  
 497 discussed. Our results have shown that the topology based on flat strips is as compact  
 498 as the classical topology based on cylindrical post resonators, while exhibiting an enhanced  
 499 broadband performance and a better power handling capability.

## 500 Open Research Section

501 Data were not used nor created for this research. The simulation results were obtained  
 502 using ANSYS High Frequency Structure Simulator (ANSYS, 2024) and CST Studio Suite  
 503 (Dassault Systemes, 2024).

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